

The Index of Refraction

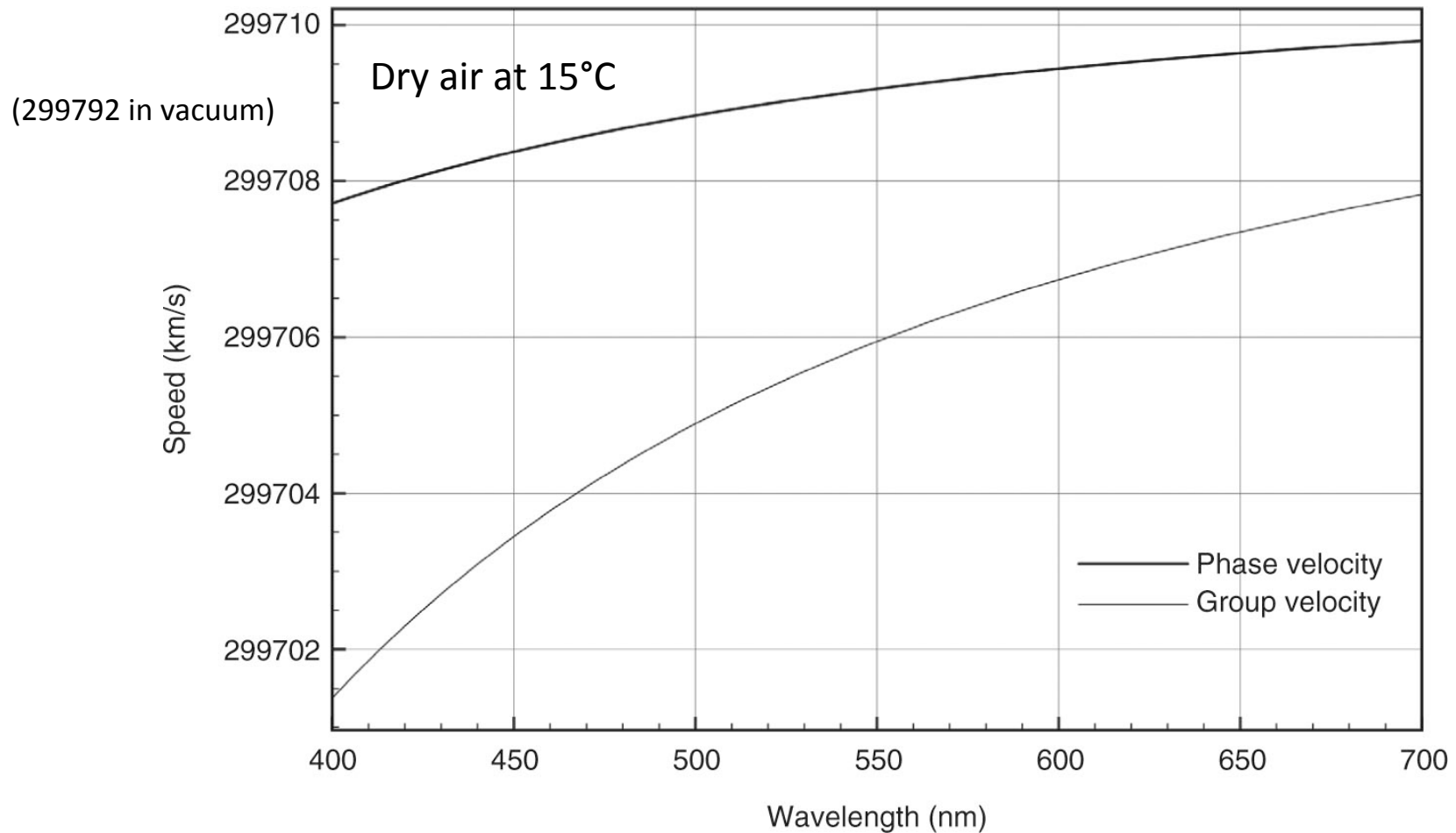
- Index of refraction is a complex quantity that is wavelength dependent

$$m(\lambda) = n(\lambda) - ik(\lambda)$$

- Real portion “n”: index of refraction
- Imaginary portion “k”: absorptivity
 - Also called “optical constants”
- Wavelength (λ) dependence is important: this is why spectroscopy works!

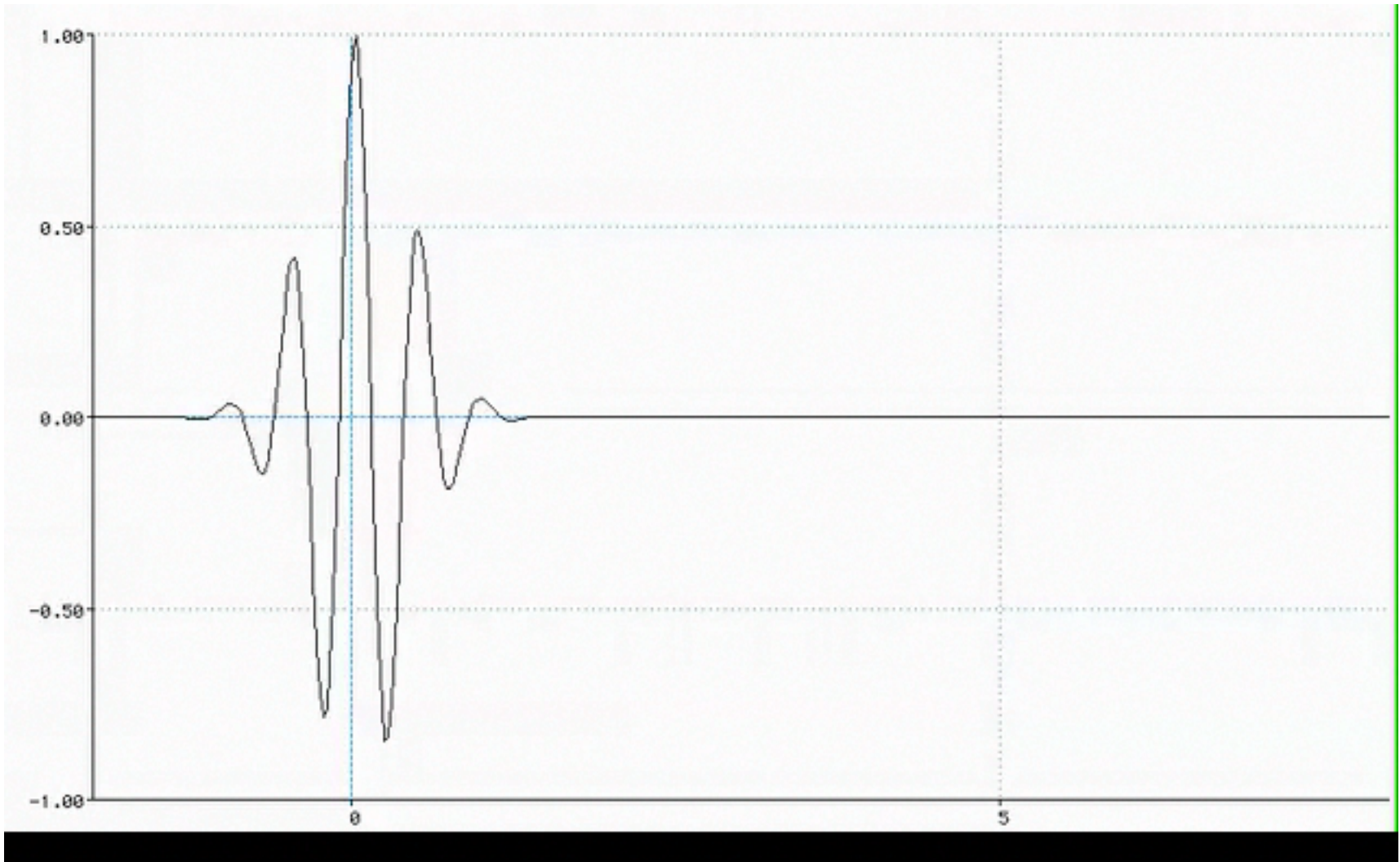
Dispersion

If n varies with wavelength, then so does phase velocity $v = \omega/k = c/n$,
... and **group velocity** $v_g = d\omega/dk$



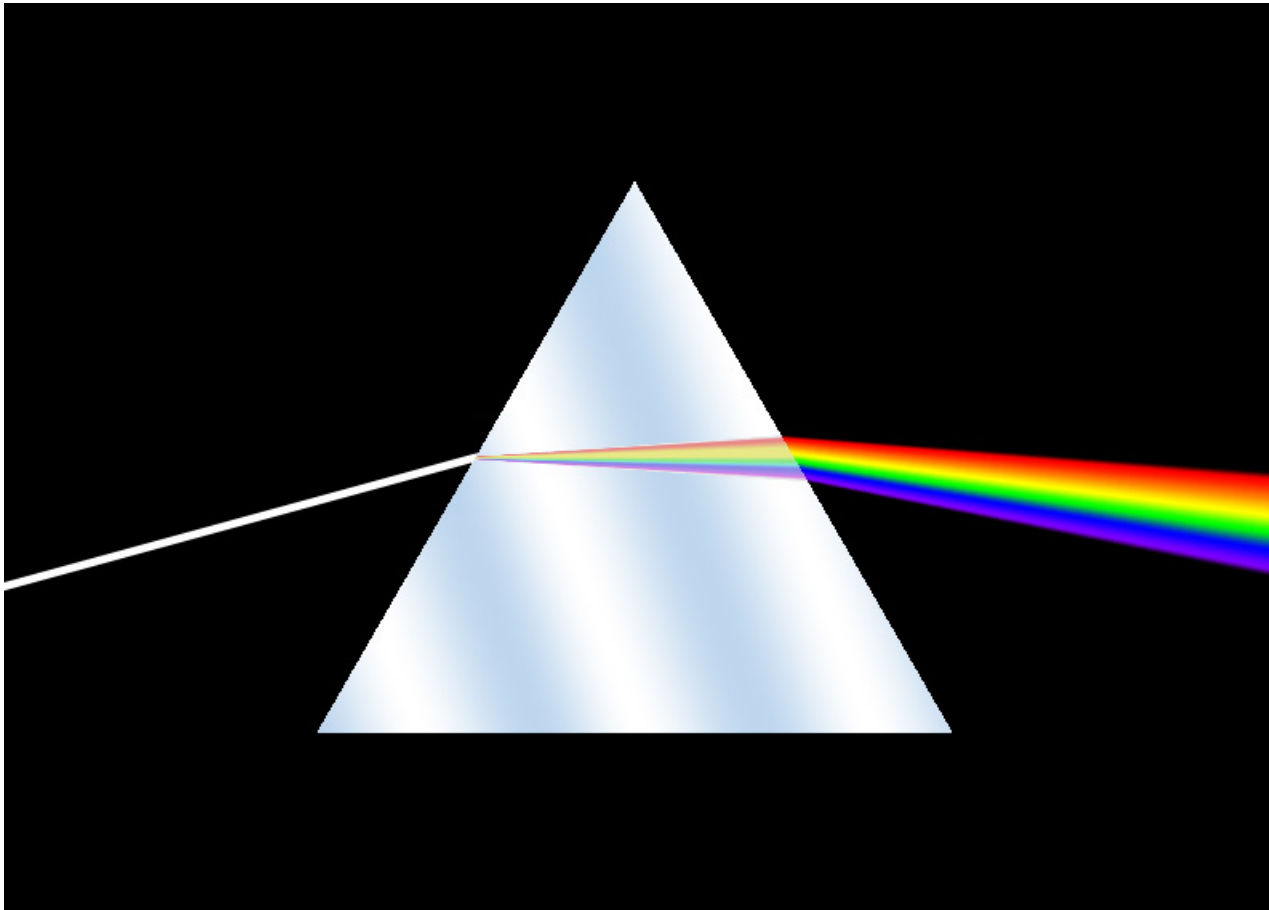
Rees Fig. 3.5

Phase velocity vs. group velocity

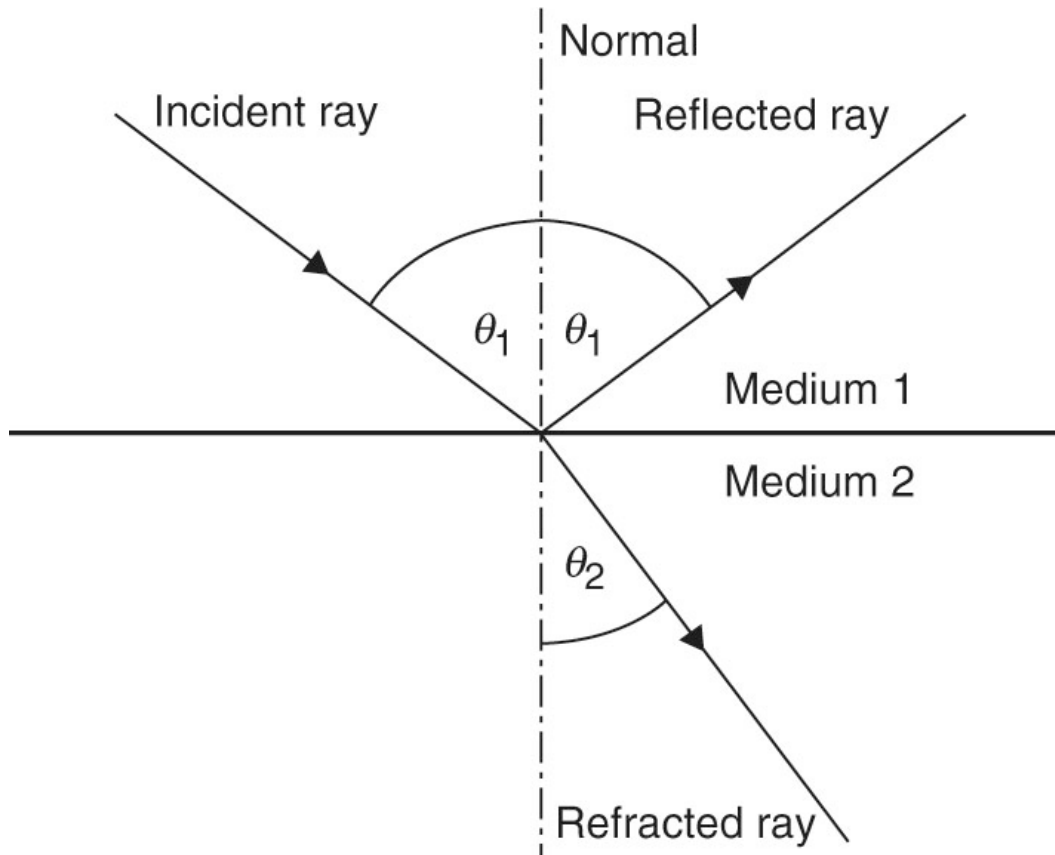


[Which is larger here?]

Refraction through a prism:
 n, k are functions of λ



Transmitted light is *refracted*

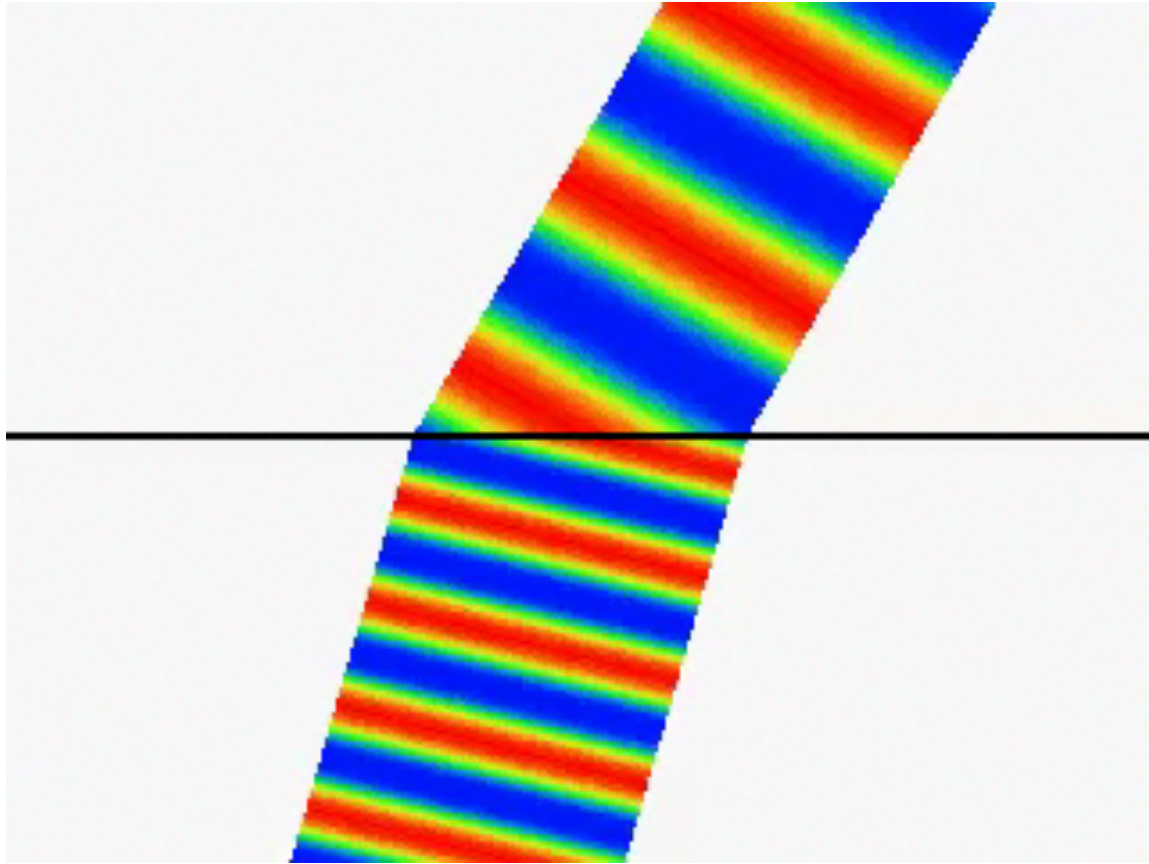


Rees Fig. 3.6

Snell's law:

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

Snell's law



Consider how refraction can distort your perception of distance...

Fresnel Reflectance

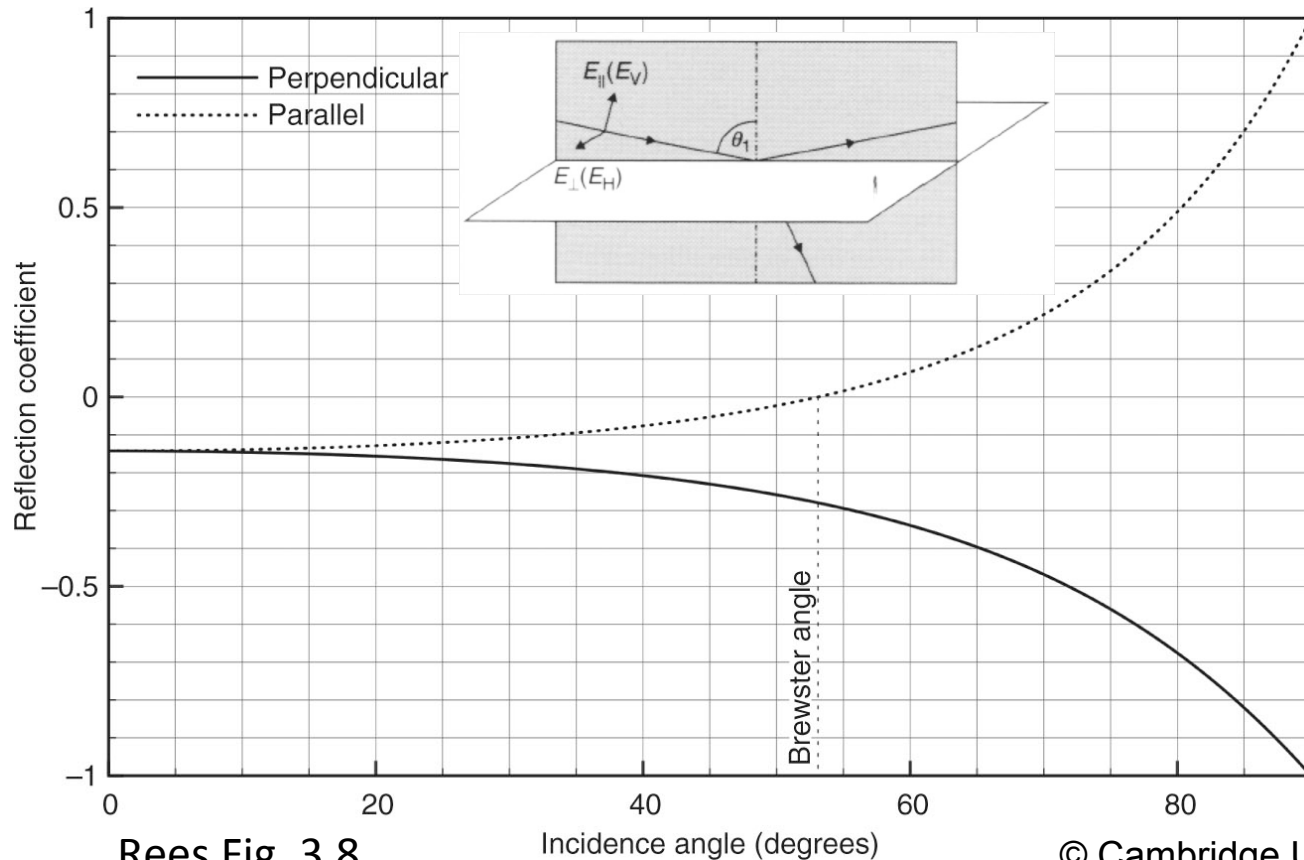
- Fresnel equations describe reflectance from a specular (smooth) surface in terms of the optical constants
 - In the special case of normal incidence:

$$R = \frac{(n - 1)^2 + k^2}{(n + 1)^2 + k^2}$$

Note: All quantities above are functions of λ .

Fresnel's equations...

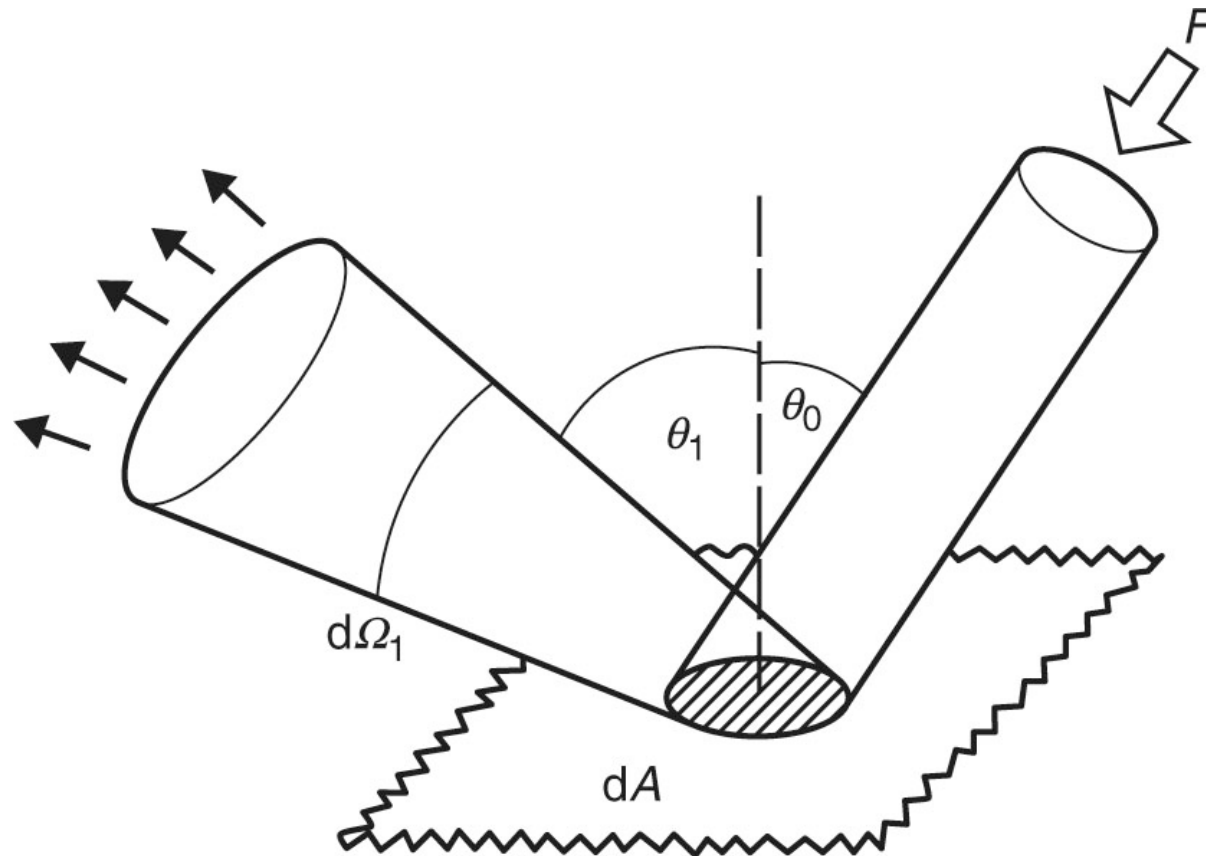
are actually much more complicated, generally giving the reflection, transmission coefficients as functions of θ_1 , θ_2 , polarization



Rees Fig. 3.8

Incidence angle (degrees)

Scattering



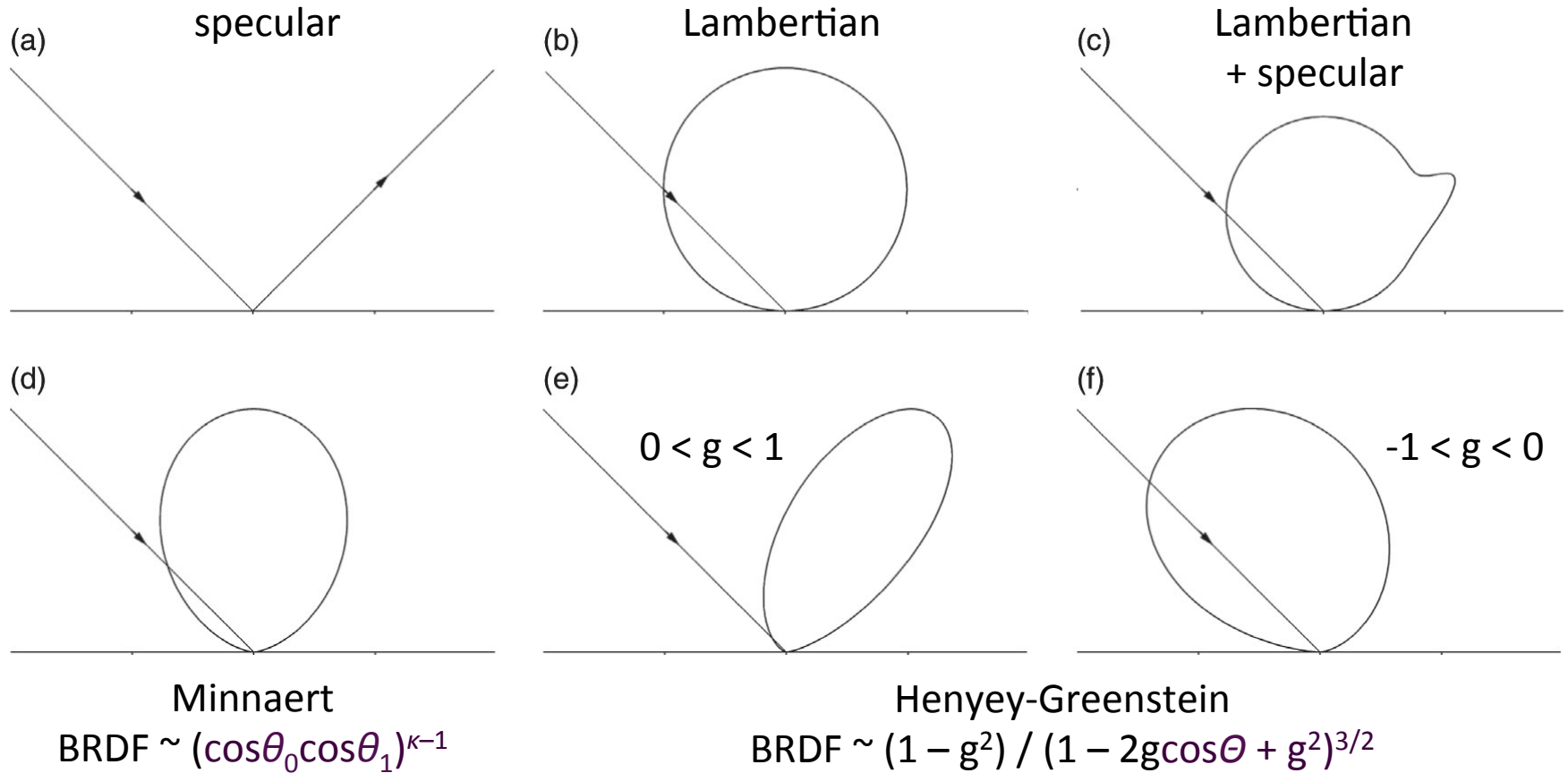
BRDF

Rees Fig. 3.9

Q: What irradiance is measured at the surface?

A: $F \cos \theta_0$

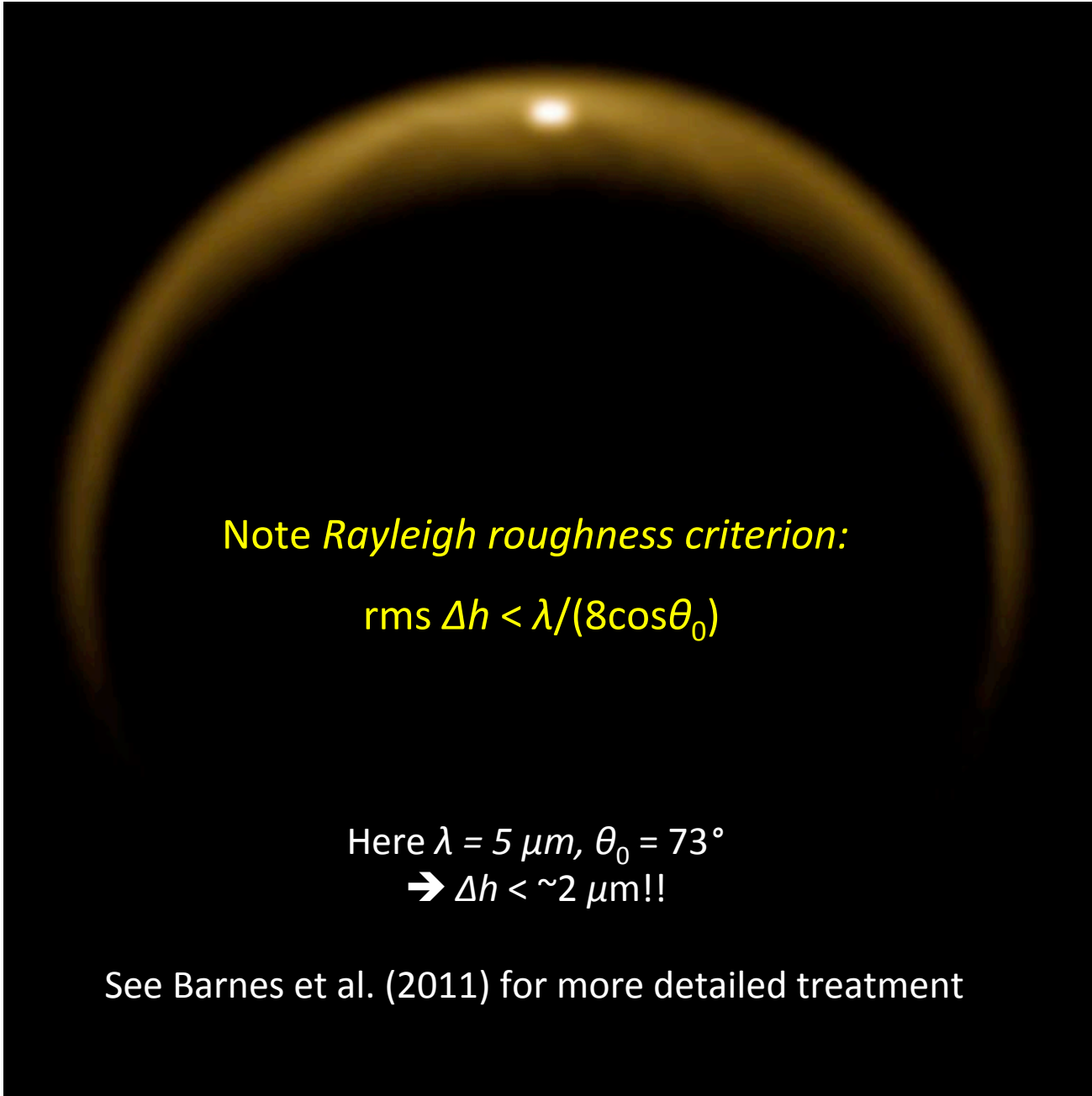
Reflection envelopes



Specular reflection: Earth



Specular reflection: Titan



Note *Rayleigh roughness criterion*:

$$\text{rms } \Delta h < \lambda / (8 \cos \theta_0)$$

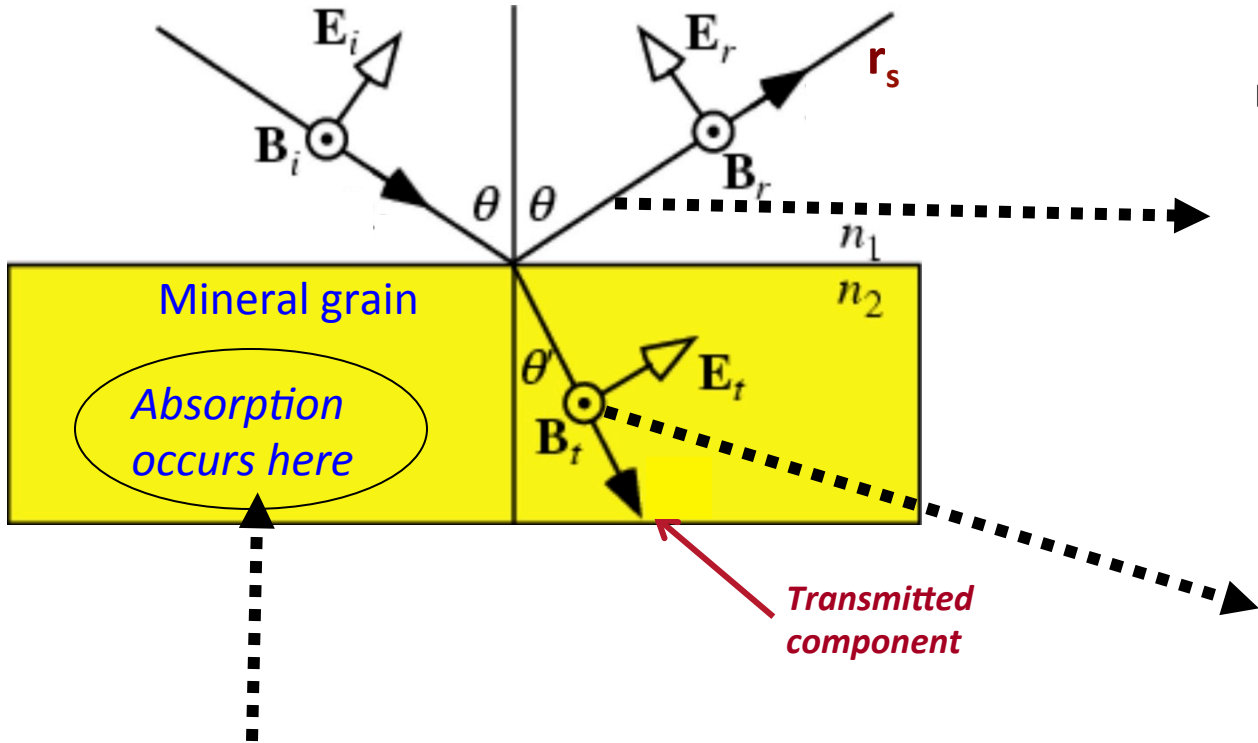
Here $\lambda = 5 \mu\text{m}$, $\theta_0 = 73^\circ$

→ $\Delta h < \sim 2 \mu\text{m}!!$

See Barnes et al. (2011) for more detailed treatment

Light is reflected, absorbed, or transmitted

The amount of specular (mirror) reflection is given by Fresnel's Law



Fresnel's law

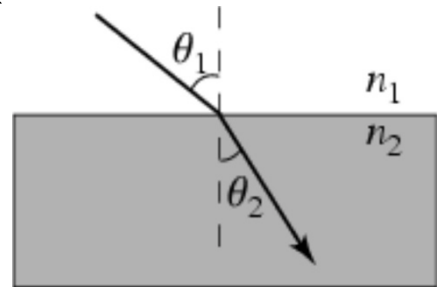
$$r_s = \frac{(n-1)^2 + K^2}{(n+1)^2 + K^2}$$

n = refractive index
 K = extinction coefficient for the solid
 r_s = fraction of light reflected from the 1st surface

Beer's law:
 (L = L₀ e^{-kz})

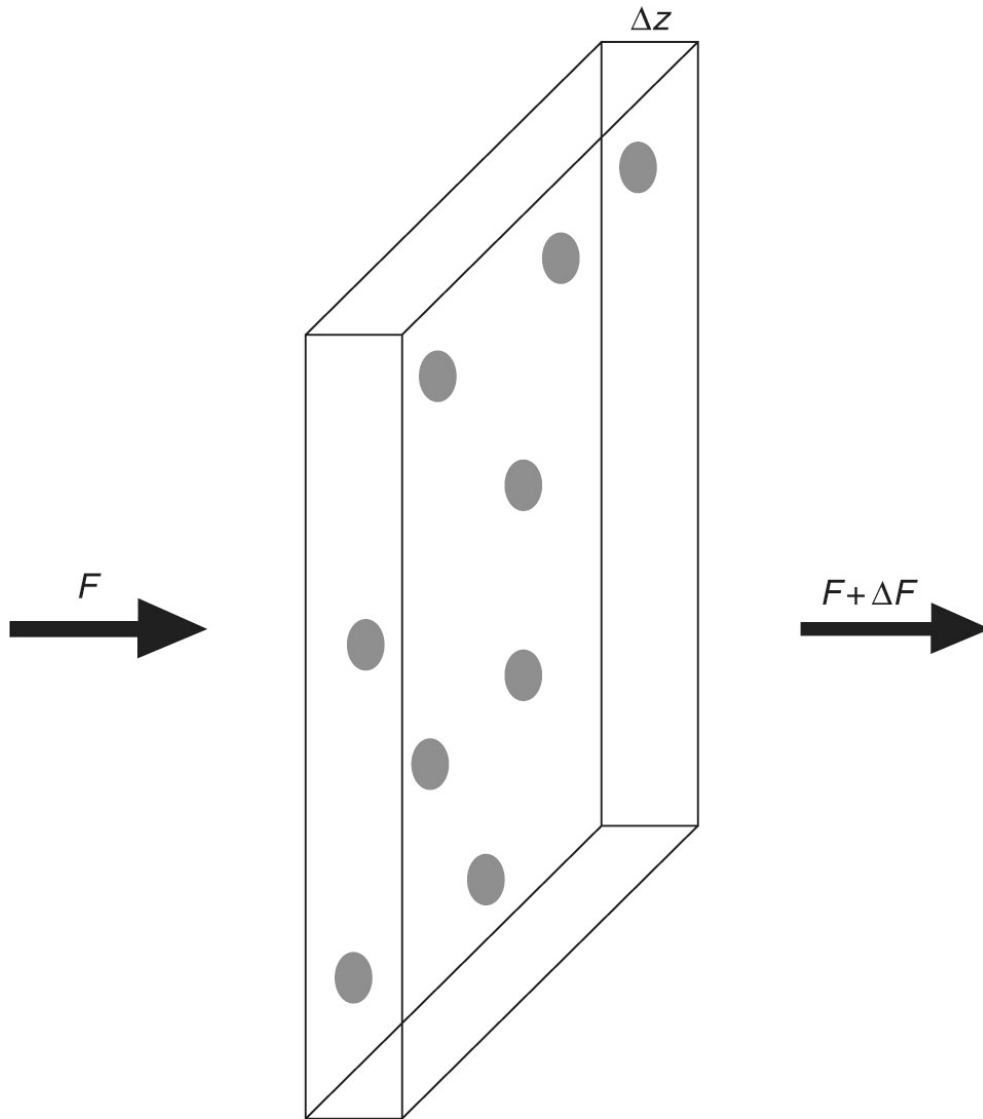
z = thickness of absorbing material
 k = absorption coefficient for the solid
 L₀ = incoming directional radiance
 L = outgoing radiance

Snell's law:
 n₁ · sinθ₁ = n₂ · sinθ₂



Light passing from one medium to another is *refracted* according to Snell's Law
n = c/v

Beer's law



$$\Delta F = -n\sigma_a F \Delta z$$

n = absorber # density

σ_a = absorption cross section

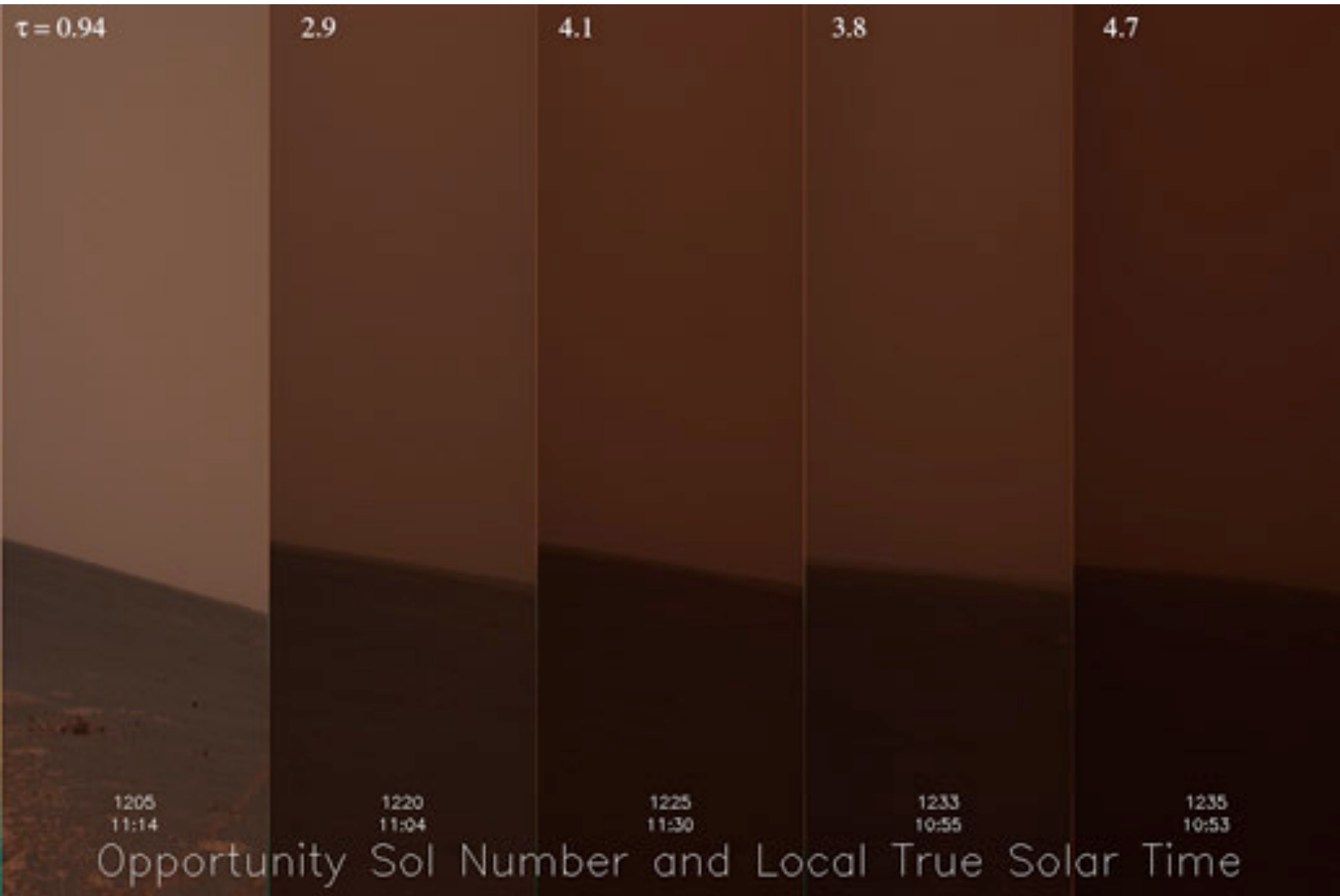
$n\sigma_a$ is the "absorption coefficient"
(1 / absorption length)

$$\rightarrow F = F_0 e^{-n\sigma_a z}$$

or $F = F_0 e^{-\tau}$
for "optical thickness" τ

Rees Fig. 3.27

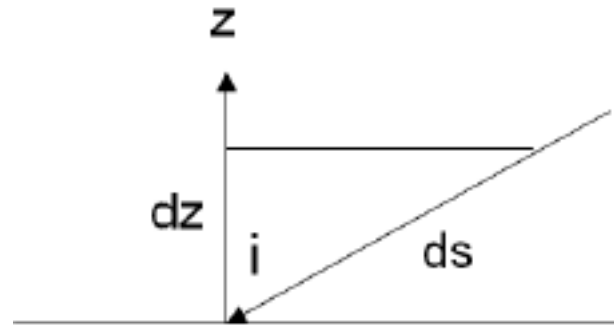
Optical thickness: the 2007 Mars dust storm



Opportunity Sol Number and Local True Solar Time

Note regarding non-normal incidence

→ *Longer path through the absorbing medium*



$$dz = ds \cos(i) = ds \mu$$

$$F = F_0 e^{-n\sigma s} = F_0 e^{-n\sigma z/\mu}$$

$$\text{i.e., } F = F_0 e^{-\tau/\mu}$$

Add emission

$$dL/dz = \gamma_a(B - L)$$

(why same coefficient on B and L?)

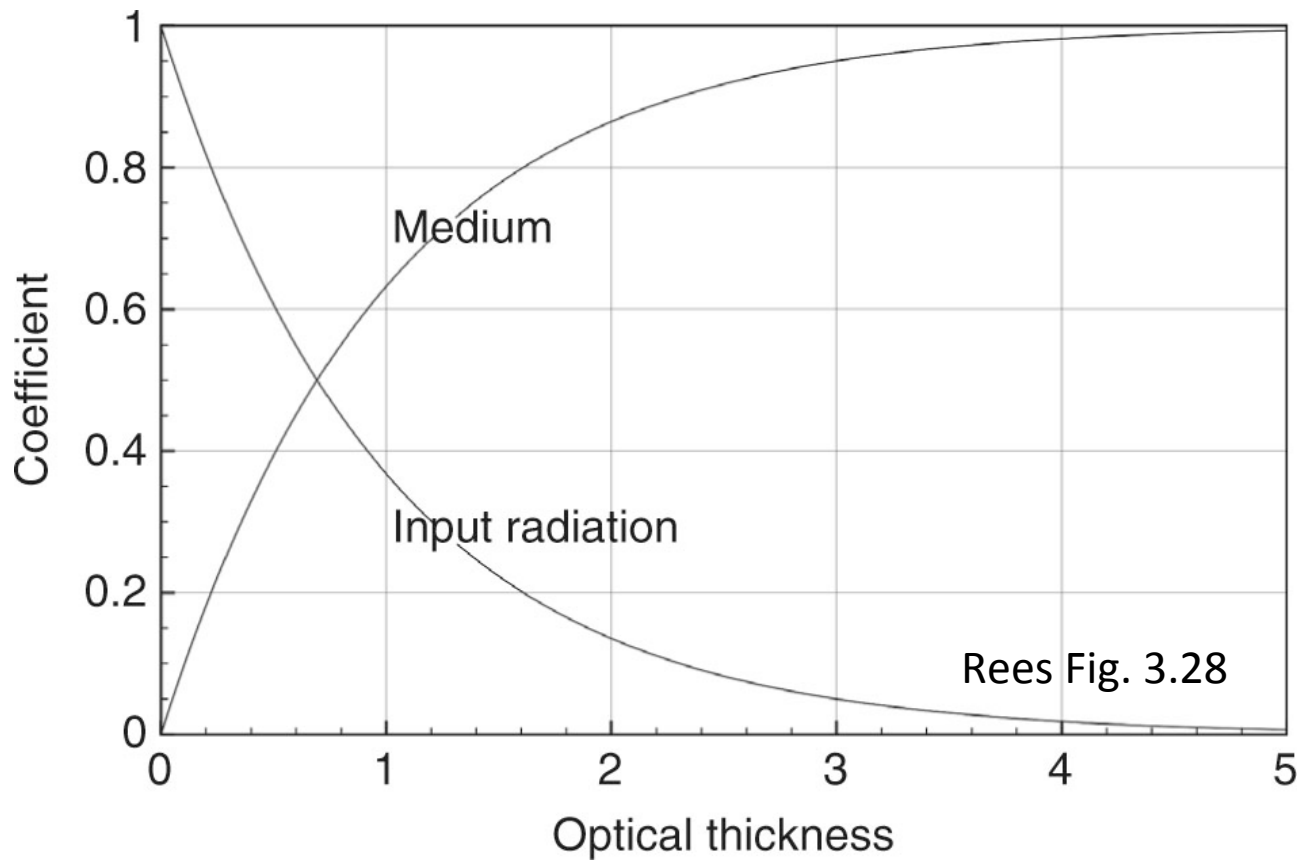
At low frequencies where $B \sim T$:

$$T_b = T_b(0)e^{-\tau} + T(1 - e^{-\tau})$$

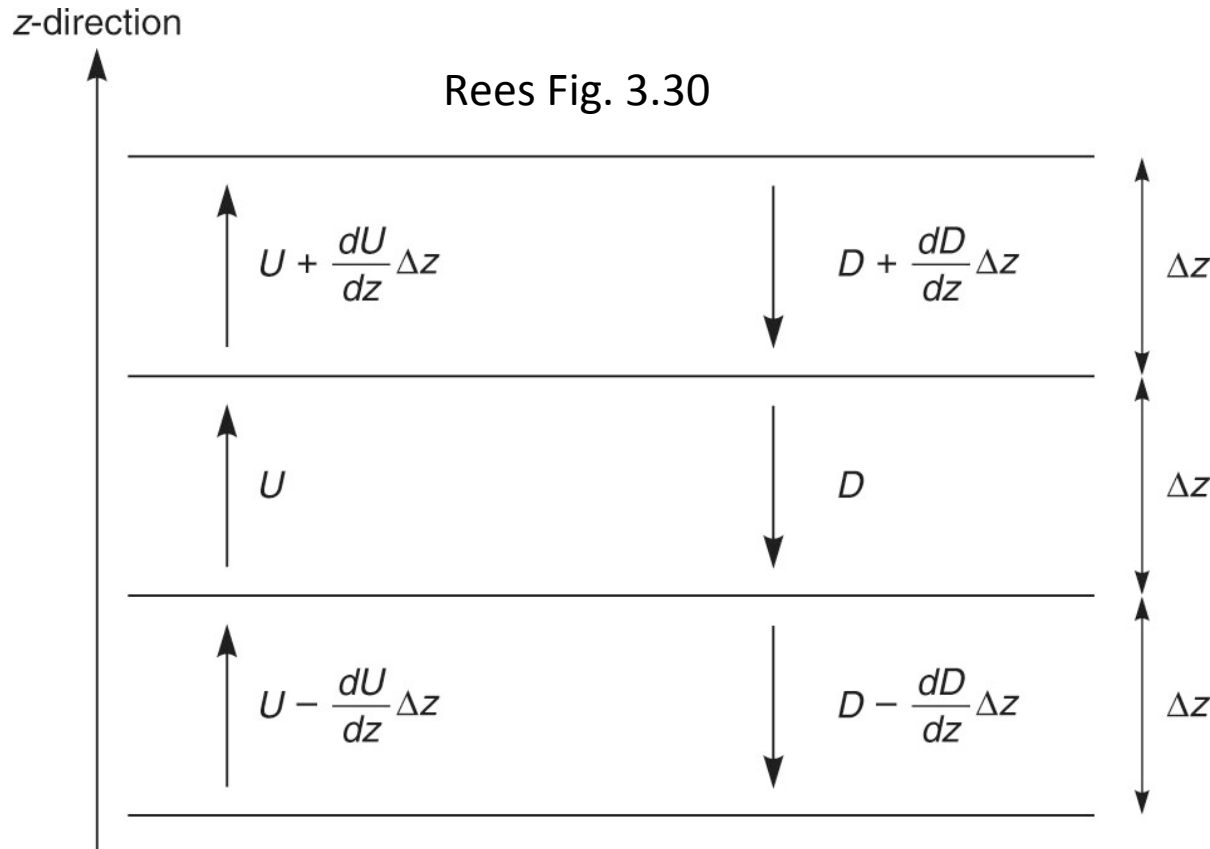
observed

input

medium



Add scattering (reflection, in 2-stream case)



$$dU/dz = -(\gamma_a + \gamma_s)U + \gamma_s D$$

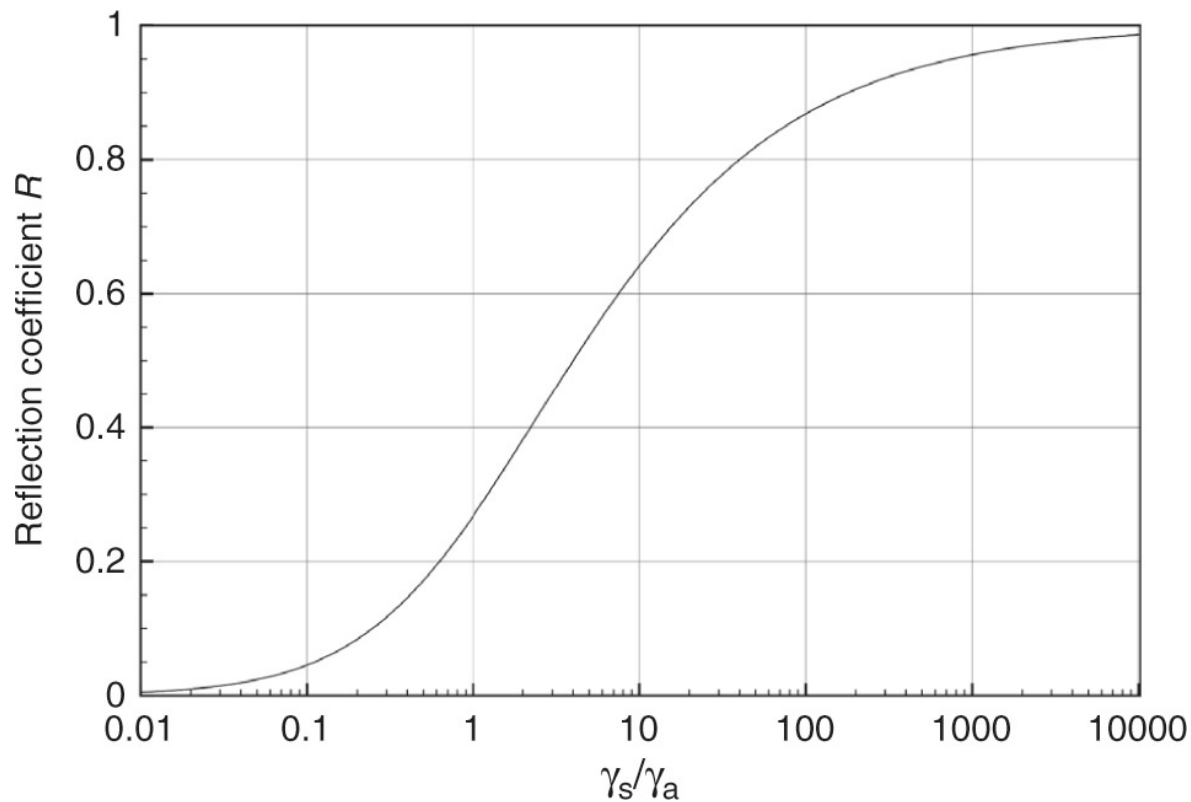
$$dD/dz = (\gamma_a + \gamma_s)D - \gamma_s U$$

Define **extinction coefficient** $\gamma_e = \gamma_a + \gamma_s$

Single scattering albedo is γ_s/γ_e

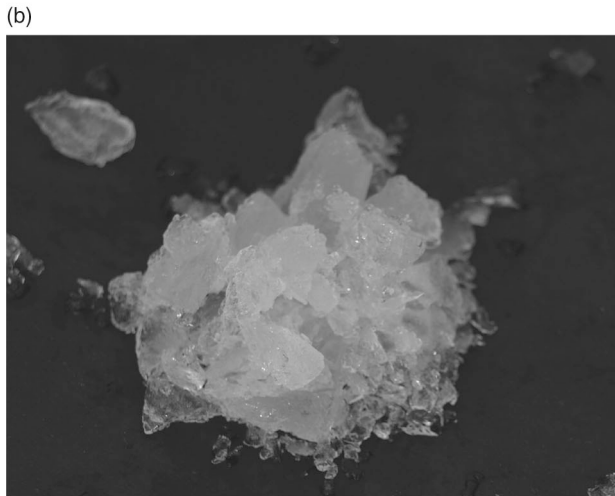
Scattering vs. absorption \rightarrow reflectance

$$R = \gamma_a/\gamma_s + 1 - \sqrt{(\gamma_a^2/\gamma_s^2 + 2\gamma_a/\gamma_s)}$$



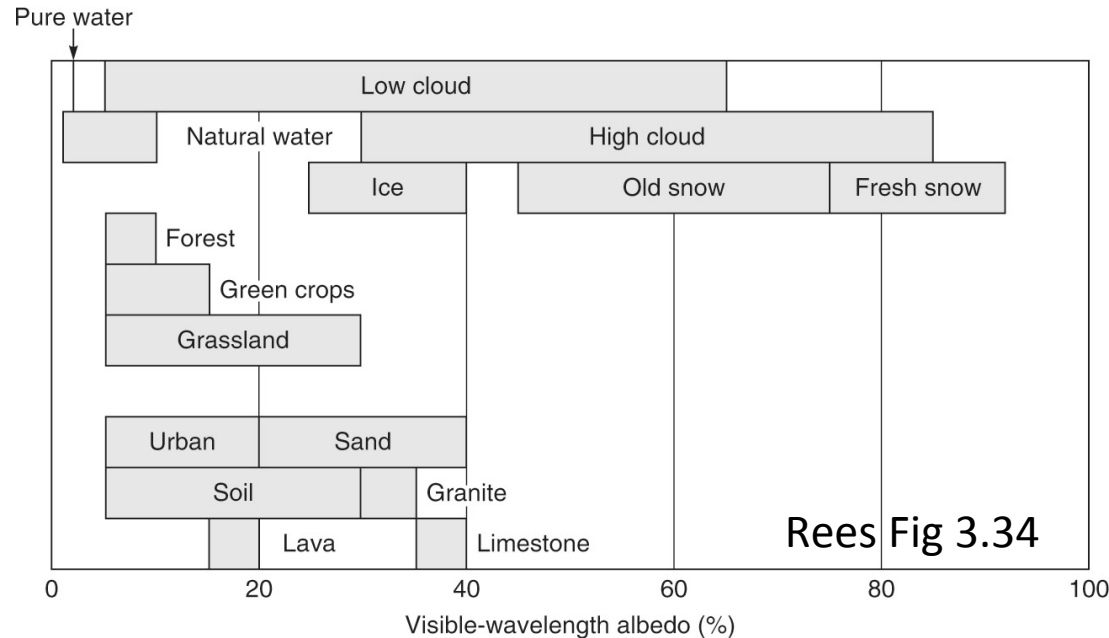
Rees Fig 3.31

Scattering vs. absorption → grain size matters



Rees Fig 3.32

Highly transmissive materials become highly reflective when you increase the density of scattering interfaces (ice, salt, clouds)



Putting it all together

$$dU/dz = -(\gamma_a + \gamma_s)U + \gamma_s D$$

$$dL/dz = \gamma_a(B - L)$$

$$\frac{dL_f(\theta, \phi)}{dz} = -(\gamma_a + \gamma_s)L_f(\theta, \phi) + \frac{\gamma_s}{4\pi} \int L_f(\theta', \phi') p(\cos\Theta) d\Omega' + \gamma_a B_f$$

$p(\cos\Theta)$ is the scattering
phase function; e.g.:

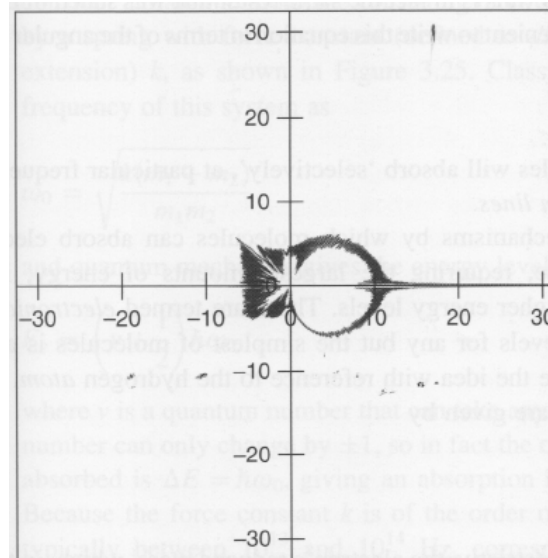


Figure 3.23. Phase function for light scattered from a spherical water droplet with a radius of 0.1 mm. The upper and lower halves of the diagram represent different polarisation states. The phase function is plotted on a logarithmic scale, and is somewhat schematic since not all details can be resolved. Note the strong forward scattering and the peak at about 140° corresponding to the primary rainbow.